

Suggested Solutions to Midterm Test for MATH4220

March 9, 2017

1. (20 points)

(a) (10 points) Find all the solutions to

$$u_x - 2u_y + 2u = 1$$

(b) (10 points) Solve the problem

$$\begin{cases} y\partial_x u + 3x^2 y\partial_y u = 0 \\ u(x=0, y) = y^2 \end{cases}$$

In which region of the xy -plane is the solution uniquely determined?

Solution:

(a) **Method 1: Coordinate Method:**

Change variables to

$$x' = x - 2y, \quad y' = -2x - y$$

Hence $u_x - 2u_y + 2u = 5u_{x'} + 2u = 1$. Thus the solution is $u(x', y') = f(y')e^{-\frac{2}{5}x'} + \frac{1}{2}$, with f an arbitrary function of one variable. Therefore, the general solutions are

$$u(x, y) = \frac{1}{2} + f(-2x - y)e^{-\frac{2}{5}(x-2y)}$$

where f is an arbitrary function.

Method 2: Geometric Method

The corresponding characteristic curves are

$$\frac{dx}{1} = \frac{dy}{-2}$$

that is, $y = -2x + C$ where C is an arbitrary constant. Then

$$\frac{d}{dx}u(x, -2x + C) = u_x(x, -2x + C) - 2u_y(x, -2x + C) = -2u(x, -2x + C) + 1$$

Hence $u(x, -2x + C) = f(C)e^{-2x} + \frac{1}{2}$, where f is an arbitrary function. Therefore,

$$u(x, y) = \frac{1}{2} + f(2x + y)e^{-2x}$$

where f is an arbitrary function.

(b) The characteristic curves are

$$\frac{dy}{3x^2 y} = \frac{dx}{y}$$

that is, $y = x^3 + C$ where C is an arbitrary constant. Then

$$\frac{d}{dx}u(x, x^3 + C) = u_x + 3x^2 u_y = 0$$

Hence $u(x, x^3 + C) = f(C)$ where f is an arbitrary function. Thus

$$u(x, y) = f(y - x^3)$$

Besides, the auxiliary condition gives that $y^2 = u(x = 0, y) = f(y)$. Hence, the solution is

$$u(x, y) = (y - x^3)^2$$

Note that when $y = 0$ the equation vanishes, thus the characteristic curves break down when $y = 0$, therefore the solution is uniquely determined on $\{(x, y) : y > 0, y > x^3\} \cup \{(x, y) : y < 0, y < x^3\} \cup \{(0, 0)\}$. (Remark: if the solution is **continuous**, then u is uniquely determined on the whole plane by the continuity of u).

2. (20 points)

(a) (4 points) What is the type of the equation

$$\partial_t^2 u + \partial_{xt}^2 u - 2\partial_x^2 u = 0 ?$$

(b) (16 points) Solve the Cauchy problem

$$\begin{cases} \partial_t^2 u + \partial_{xt}^2 u - 2\partial_x^2 u = 2, & -\infty < x < +\infty, \quad -\infty < t < +\infty \\ u(x, t = 0) = x^2, \quad \partial_t u(x, t = 0) = 0, & -\infty < x < +\infty \end{cases}$$

Solution:

(a) Since $a_{11} = 1, a_{12} = \frac{1}{2}, a_{22} = -2$, then $a_{12}^2 - a_{11}a_{22} = \frac{9}{4} > 0$, hence it is hyperbolic.

(b) Let

$$t = t', \quad x = \frac{1}{2}t' + \frac{3}{2}x'$$

and $v(x', t') = u(x, t)$, then v satisfies

$$\begin{cases} \partial_{t'}^2 v - \partial_{x'}^2 v = 2, \\ v(x', t' = 0) = (\frac{3}{2}x')^2, \quad \partial_{t'} v(x', t' = 0) = \frac{3}{2}x', \end{cases}$$

Thus d'Alembert formula gives that

$$\begin{aligned} v(x', t') &= \frac{1}{2} \left\{ \frac{9}{4}(x' + t')^2 + \frac{9}{4}(x' - t')^2 \right\} + \frac{1}{2} \int_{x'-t'}^{x'+t'} \frac{3}{2} y dy + \frac{1}{2} \int_0^{t'} \int_{x'-(t'-s)}^{x'+(t'-s)} 2 dy ds \\ &= \frac{13}{4}t'^2 + \frac{9}{4}x'^2 + \frac{3}{2}x't' \end{aligned}$$

Then

$$u(x, t) = v(x', t') = v\left(\frac{2}{3}x - \frac{1}{3}t, t\right) = 3t^2 + x^2.$$

3. (20 points)

(a) (5 points) State the definition of a well-posed PDE problem.

(b) (5 points) Is the following problem well-posed? Why?

$$\begin{cases} \frac{d^2 u}{dx^2} + \frac{du}{dx} = 1, & 0 < x < 1 \\ u'(0) = 1, u'(1) = 0 \end{cases}$$

- (c) (10 points) State and prove the uniqueness and continuous dependence of solutions to the problem

$$\begin{cases} \partial_t u = \partial_x^2 u, & 0 < x < 1, \quad 0 < t < T, \quad T > 0 \\ \partial_x u(0, t) = 0, \partial_x u(1, t) = 0, & t > 0 \\ u(x, t = 0) = \phi(x), & 0 \leq x \leq 1 \end{cases}$$

Solution:

- (a) A PDE problem is said to be well-posed if the following three properties are satisfied:

Existence: There exists at least one solution $u(x, t)$ satisfying all these conditions.

Uniqueness: There is at most one solution.

Stability: The unique solution $u(x, t)$ depends in a stable manner on the data of the problem. This means that if the data are changed a little, the corresponding solution changes only a little.

- (b) The problem is **not** well-posed, since the solution doesn't exist.

Indeed, the general solution to the ODE $\frac{d^2 u}{dx^2} + \frac{du}{dx} = 1$ is $u(x) = C_1 + C_2 e^{-x} + x$. Then $u'(x) = -C_2 e^{-x} + 1$, and boundary condition $u'(0) = 1$ implies that $C_2 = 0$, however $u'(1) = 1 \neq 0$. Thus the solution cannot exist.

- (c) **Uniqueness:** Let u_1 and u_2 be any two solutions to the problem, then $u_1 = u_2$.

Continuous dependence on initial data: If u_1 and u_2 are solutions to the problem with initial condition $\phi_1(x)$ and $\phi_2(x)$ respectively, then

$$\sup_{0 \leq t \leq T} \int_0^1 |u_1 - u_2|^2 dx \leq \int_0^1 |\phi_1 - \phi_2|^2 dx.$$

Proof:

(Uniqueness): Let $v = u_1 - u_2$, then v satisfies

$$\begin{cases} \partial_t v = \partial_x^2 v, & 0 < x < 1, \quad 0 < t < T, \quad T > 0 \\ \partial_x v(0, t) = 0, \partial_x v(1, t) = 0, & t > 0 \\ v(x, t = 0) = 0, & 0 \leq x \leq 1 \end{cases}$$

Multiplying the both sides of $\partial_t v = \partial_x^2 v$ by v and taking intergral from 0 to 1 with respect to x , then we have

$$\int_0^1 \partial_t v v dx = \int_0^1 \partial_x^2 v v dx$$

Then

$$L.H.S = \frac{d}{dt} \int_0^1 \frac{1}{2} v^2 dx$$

$$R.H.S = \partial_x v v \Big|_0^1 - \int_0^1 (\partial_x v)^2 dx = - \int_0^1 (\partial_x v)^2 dx \leq 0$$

Then, we have for $t > 0$

$$0 \leq \int_0^1 \frac{1}{2} v^2(x, t) dx \leq \int_0^1 \frac{1}{2} v^2(x, 0) dx = 0$$

By the continuity of v , we have $v(x, t) \equiv 0$, $0 < x < 1, 0 < t < T$. Thus we have shown that $u_1(x, t) \equiv u_2(x, t)$ for $0 < x < 1, 0 < t < T$.

(Continuous dependence on initial data:) Let $v = u_1 - u_2$, then v satisfies

$$\begin{cases} \partial_t v = \partial_x^2 v, & 0 < x < 1, \quad 0 < t < T, \quad T > 0 \\ \partial_x v(0, t) = 0, \partial_x v(1, t) = 0, & t > 0 \\ v(x, t = 0) = \phi_1(x) - \phi_2(x) =: \tilde{\phi}(x), & 0 \leq x \leq 1 \end{cases}$$

Multiplying the both sides of $\partial_t v = \partial_x^2 v$ by v and taking intergral from 0 to 1 with respect to x , then we have

$$\int_0^1 \partial_t v v dx = \int_0^1 \partial_x^2 v v dx$$

Then

$$L.H.S = \frac{d}{dt} \int_0^1 \frac{1}{2} v^2 dx$$

$$R.H.S = \partial_x v v \Big|_0^1 - \int_0^1 (\partial_x v)^2 dx = - \int_0^1 (\partial_x v)^2 dx \leq 0$$

Then, we have

$$\sup_{0 \leq t \leq T} \int_0^1 \frac{1}{2} v^2(x, t) dx \leq \int_0^1 \frac{1}{2} v^2(x, 0) dx = \int_0^1 \frac{1}{2} \tilde{\phi}^2(x) dx$$

which completes the proof of the continuous dependence on initial data.

4. (20 points)

(a) (15 points) Derive the solution formula for the following initial-boundary value problem

$$\begin{cases} \partial_t u = \partial_x^2 u, & 0 < x < +\infty, \quad t > 0 \\ u(x, t = 0) = \phi(x) & 0 < x < +\infty \\ \partial_x u(x = 0, t) = 0, & t > 0 \end{cases}$$

by the method of reflection (with all the details of the derivation).

(b) (5 points) Let $\phi(x) = \cos x, 0 < x < +\infty$. Find the maximum value of $u(x, t)$.

Solution:

(a) Use the reflection method, and first consider the following Cauchy Problem:

$$\begin{cases} \partial_t v = \partial_x^2 v, & 0 < x < +\infty, \quad t > 0 \\ v(x, t = 0) = \phi_{even}(x) & 0 < x < +\infty \end{cases}$$

where $\phi_{even}(x)$ is even extension of ϕ which is given by

$$\phi_{even}(x) = \begin{cases} \phi(x), & \text{if } x > 0 \\ \phi(-x), & \text{if } x < 0 \end{cases}$$

Then the unique solution is given by:

$$v(x, t) = \int_{-\infty}^{\infty} S(x - y, t) \phi_{even}(y) dy$$

And since $\phi_{even}(x)$ is even, so is $v(x, t)$ for $t > 0$, which implies

$$\partial_x v(x = 0, t) = 0, t > 0$$

Set $u(x, t) = v(x, t), x > 0$, then $u(x, t)$ is the unique solution of Neumann Problem on the half-line. More presicely, $x > 0, t > 0$

$$\begin{aligned} u(x, t) &= \int_0^{\infty} S(x - y, t) \phi(y) dy + \int_{-\infty}^0 S(x - y, t) \phi(-y) dy \\ &= \int_0^{\infty} S(x - y, t) \phi(y) dy + \int_0^{\infty} S(x + y, t) \phi(y) dy \\ &= \frac{1}{\sqrt{4k\pi t}} \int_0^{\infty} [e^{-\frac{(x-y)^2}{4kt}} + e^{-\frac{(x+y)^2}{4kt}}] \phi(y) dy. \end{aligned}$$

Here, $k = 1$.

(b) By (a), the solution is given by ($k = 1$)

$$u(x, t) = \frac{1}{\sqrt{4k\pi t}} \int_0^\infty [e^{-\frac{(x-y)^2}{4kt}} + e^{-\frac{(x+y)^2}{4kt}}] \cos y dy$$

then

$$|u(x, t)| \leq \frac{1}{\sqrt{4k\pi t}} \int_0^\infty [e^{-\frac{(x-y)^2}{4kt}} + e^{-\frac{(x+y)^2}{4kt}}] |\cos y| dy \leq \frac{1}{\sqrt{4k\pi t}} \int_0^\infty [e^{-\frac{(x-y)^2}{4kt}} + e^{-\frac{(x+y)^2}{4kt}}] dy = 1$$

That is, $|u(x, t)| \leq 1$. Note that $\max_{0 < x < \infty} u(x, 0) = \max_{0 < x < \infty} \cos x = 1$, which implies that 1 can be attained by u . Hence

$$\max_{0 < x < \infty, t > 0} u(x, t) = 1.$$

5. (20 points)

(a) (10 points) Prove the following generalized maximum principle:

If $\partial_t u - k\partial_x^2 u \leq 0$ on $R \triangleq [0, l] \times [0, T]$ with a positive constant k , then

$$\max_R u(x, t) = \max_{\partial R} u(x, t)$$

here $\partial R = \{(x, t) \in R \mid \text{either } t = 0, \text{ or } x = 0, \text{ or } x = l\}$.

(b) (10 points) Show if $v(x, t)$ solves the following problem

$$\begin{cases} \partial_t v = k\partial_x^2 v + f(x, t), & 0 < x < l, 0 < t < T \\ v(x, 0) = 0, & 0 < x < l \\ v(0, t) = 0, v(l, t) = 0, & 0 \leq t \leq T \end{cases}$$

with a continuous function f on $R \triangleq [0, l] \times [0, T]$. Then

$$v(x, t) \leq t \max_R |f(x, t)|$$

(Hint, consider $u(x, t) = v(x, t) - t \max_R |f(x, t)|$ and apply the result in (a).)

Solution:

(a) Let $v(x, t) = u(x, t) + \epsilon x^2$, then v satisfies

$$\partial_t v - k\partial_x^2 v = \partial_t u - k\partial_x^2 u - 2k\epsilon < 0$$

First, **claim** that v attains its maximum on the parabolic boundary R . Let $\max_R v(x, t) = M = v(x_0, t_0)$. Suppose on the contrary, then either

i. $0 < x_0 < l, 0 < t_0 < T$.

In this case, $v_t(x_0, t_0) = v_x(x_0, t_0) = 0$ and $v_{xx}(x_0, t_0) \leq 0$. Thus $\partial_t v - k\partial_x^2 v|_{(x_0, t_0)} \geq 0$, which is impossible.

ii. $0 < x_0 < l, t_0 = T$.

In this case, $v_t(x_0, t_0) \geq 0, v_x(x_0, t_0) = 0$ and $v_{xx}(x_0, t_0) \leq 0$. Thus $\partial_t v - k\partial_x^2 v|_{(x_0, t_0)} \geq 0$, which is impossible.

Hence

$$\max_R v(x, t) = \max_{\partial R} v(x, t).$$

Then for any $(x, t) \in R$,

$$u(x, t) \leq v(x, t) + \epsilon x^2 \leq \max_{\partial R} v(x, t) \leq \max_{\partial R} u(x, t) + \epsilon l^2$$

Letting $\epsilon \rightarrow 0$ gives $u(x, t) \leq \max_{\partial R} u(x, t)$ for any $(x, t) \in R$. Hence $\max_R u(x, t) = \max_{\partial R} u(x, t)$.

(b) Let $u(x, t) = v(x, t) - t \max_R |f(x, t)|$, then u satisfies

$$\begin{cases} \partial_t u - k \partial_x^2 u = -\max_R |f(x, t)| + f(x, t) \leq 0 \\ u(x, 0) = 0, \\ u(0, t) = u(l, t) = -t \max_R |f(x, t)| \leq 0 \end{cases}$$

Hence the result in (a) implies that for any $(x, t) \in R$,

$$u(x, t) \leq \max_{\partial R} u(x, t) = 0$$

that is, $v(x, t) \leq t \max_R |f(x, t)|$.